## STRUCTURE FUNCTIONS OF STAR-BRANCHED POLYELECTROLYTES

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Unusual structure functions have been measured for salt-free aqueous solutions of sodium sulfonated polystyrene (NaPSS) star-branched polyelectrolytes. Whatever the concentration, they display two kinds of maxima (figure 1). The first is related to a position order between star cores. It covers a main maximum at  $q_1^*$  and a further harmonic at  $\sqrt{3} q_1^*$  for the lower concentrations in the dilute regime. The second comes into sight at higher q values. It involves one maximum at  $q_2^*$ , quite similar to the broad halo observed in the scattering pattern of a semidilute aqueous solution of NaPSS linear polyelectrolytes. It is therefore associated with a correlation hole of electrostatic character.

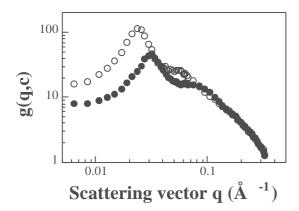
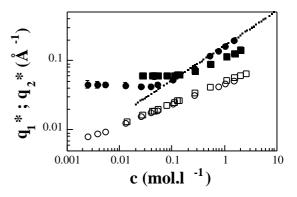


Figure 1: Structure functions of NaPSS star solutions obtained on the spectrometer PACE at LLB The NaPSS star sample corresponds to f=12,  $N_a=100$  and  $t_S=1$ . Two concentrations are considered:  $c=0.102 \text{ mol.}\Gamma^1$  (open circles);  $0.258 \text{ mol.}\Gamma^1$  (dark circles)

The variations of  $q_1^*$  and  $q_2^*$  with the monomer concentration c allow to stress the difference in nature of these maxima (Figure 2).

 $q_1^*$  almost scales as  $c^{1/3}$  with prefactors depending on the star functionality f and the arm degree of polymerization  $N_a$ . No dependence upon the arm degree of sulfonation  $\tau_S$ , has been found in the  $\tau_S$ -range providing solubility of the NaPSS stars. Obviously, such results support the close relation between  $q_1^*$  and the mean distance between star cores  $(d = 2\pi/q_1^*)$ , i.e.

the short or long range order character of the first maximum.



**Figure 2**: Scattering vectors  $q_1^*$  and  $q_2^*$  associated with the main maxima of the NaPSS star structure functions Two NaPSS star samples are considered. They have the same functionality (f=12) and arm degree of polymerization ( $N_a$  =100), but distinct  $\mathbf{t}_S$  values. Open circles and squares correspond to  $q_1^*$  values measured for  $\mathbf{t}_S$ =1 and 0.65, respectively. Dark circles and squares correspond to  $q_2^*$  values measured for  $\mathbf{t}_S$ =1 and 0.65, respectively. The scaling law related to the electrostatic peak of NaPSS linear polyelectrolytes in semidilute aqueous solutions is also reported for  $\mathbf{t}_S$ =1 (dotted line)

The c-dependence of  $q_2^*$  is quite different  $(q_1^*$  and  $q_2^*$ are therefore not commensurate). In the dilute regime  $(c < c^*)$ ,  $q_2^*$  is found to be constant, only depending on  $N_a$  and  $\tau_S$ . On the other hand, in the semidilute regime (c>c\*),  $q_2$ \* scales as  $c^{\alpha}$  with  $\alpha$  in between 1/3 and 1/2, according to the  $\tau_S$  value. Moreover, it is unaffected by any change in N<sub>a</sub>. Such scaling laws fit those associated with the position q\* of the broad peak characterizing the structure functions of NaPSS linear polyelectrolytes in semidilute aqueous solutions. The second maximum of the NaPSS star structure functions is therefore primarily due to the existence of a correlation hole (or tube) surrounding each polyion, from which the others are expelled by the electrostatic repulsion. For c>c\*, it mainly corresponds to the semidilute solution formed by the interpenetration of stars and q<sub>2</sub>\* yields its mean mesh size  $\xi$  or the Debye-Hückel screening length  $\kappa^{-1}$ (isotropic model). For c<c\*, it still represents the

electrostatic repulsion between arms. However, in that case, it is exclusively concerned with arms belonging to the same star. It is indeed a simple intramolecular feature, as demonstrated by using the zero average contrast method. Then  $q_2^*$  yields an effective screening length  $\kappa^{-1}$ , through the relation  $q_2^*\!\!\propto\!\!\kappa$ , which gives an information about the condensation of couterions inside each star. From this point of view, it is not really surprising that  $q_2^*$  is almost constant in the dilute regime. However, the results presented in figure 2 also

show that the counterion condensation inside each star increases as  $\tau_{S}$  decreases. On the other hand, the variation of  $q_{2}^{*}$  with  $N_{a}$  is in agreement with the existence of an internal screening length distribution. Finally, the distinct behaviours of  $q_{2}^{*}$  in both the dilute and semidilute regimes provide a way to estimate the overlap concentration  $c^{*}.$